

1                   **YOUNG CLOSE-BY NEUTRON STARS: THE GOULD BELT**  
 2                   **VS. THE GALACTIC DISC**

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9                   (Received 05 September 2004; accepted 18 October 2004)

10 **Abstract.** We present new population synthesis calculations of close young neutron stars. In compar-  
 11 ison with our previous investigation we use a different neutron star mass spectrum and different initial  
 12 spatial and velocity distributions. The results confirm that most of ROSAT dim radioquiet isolated  
 13 neutron stars had their origin in the Gould Belt. We predict that about several tens of young neutron  
 14 stars can be identified in ROSAT All Sky Survey data at low galactic latitudes. Some of these sources  
 15 also can have counterparts among EGRET unidentified sources.

16 **Keywords:** stars: neutron, stars: evolution, stars: statistics, X-ray: stars

17                   **1. Introduction**

18 Over the last decade X-ray missions revealed an increasing number of isolated  
 19 neutron stars (INs) in the solar vicinity. Many of these sources, essentially discov-  
 20 ered by ROSAT, are not observed as active radio pulsars and show quite peculiar  
 21 emission properties, both at X-ray and optical wavelengths (see, e.g. Treves et al.,  
 22 2000; Becker and Pavlov, 2002; Haberl, 2004 for recent reviews). Their spectrum  
 23 is peaked at  $\sim 100$  eV and is well described in terms of a featureless blackbody.  
 24 The optical emission (when observed, see e.g. Kaplan et al., 2003 on the case of  
 25 RX J0720-3125) appears close to a Rayleigh-Jeans distribution but lies well above  
 26 the extrapolation of the X-ray blackbody to optical wavelengths.

27 The many puzzling features of X-ray emitting INs offer contrasted views about  
 28 their nature. Although several interpretations have been proposed (in terms of  
 29 old INs accreting the interstellar medium, decaying magnetars or even quark  
 30 stars), the more conservative explanation is that they are conventional middle-aged  
 31 ( $\approx 10^5$ – $10^6$  years) cooling NSs which for reasons not understood as yet fail to be  
 32 detected as radio emitters.

33 A possible problem with this latter scenario is connected with the observed  
 34 overabundance of these sources in the solar proximity with respect to what predicted  
 35 by population synthesis models. If INs are born in the galactic disc (at about the



*Astrophysics and Space Science xxx: 1–11, 2004.*

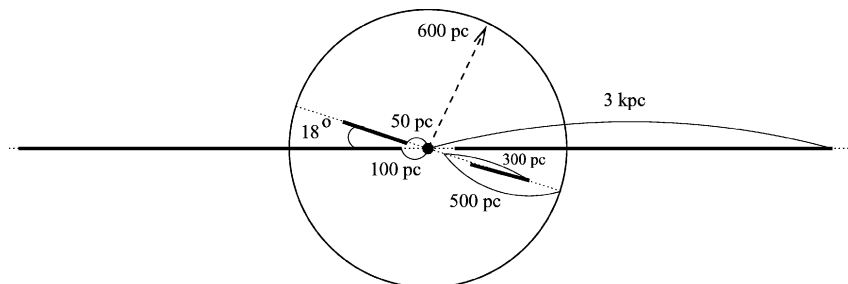
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solar distance from the galactic center) at the same rate at which radio pulsars are 36  
 formed and if they follow a standard cooling history then the number of detectable 37  
 X-ray sources falls short of the observed one by about a factor a few (Neuhäuser 38  
 and Trümper, 1999; Popov et al., 2000). A possible solution is to invoke a recent 39  
 epoch of enhanced NS formation in  $\lesssim 1$  kpc around the Sun. Originally this idea has 40  
 been suggested by Grenier (2000) and Gehrels et al. (2000) in connection with the 41  
 possibility that unidentified EGRET sources are young close-by NSs. The Gould 42  
 Belt, a collection of young star associations which encompasses the Sun, appears 43  
 the most likely birthplace for the majority of these NSs (see the Belt description in 44  
 Pöppel, 1997). In Popov et al. (2002) it was suggested that INSs observed by ROSAT 45  
 as dim X-ray sources can be explained as young cooling objects originated mainly 46  
 from the Gould Belt. Very recently Popov et al. (2003) (hereafter Paper I) addressed 47  
 this issue in detail by means of a population synthesis model in which NS formation 48  
 in the Belt (in addition to that in the galactic disc) was properly accounted for. 49

In this paper we present some refinements to the results of Paper I. In particular 50  
 we explore the effects of modifying and relaxing some of the original assumptions 51  
 contained in Paper I on the computed  $\log N - \log S$  of cooling NSs. A central 52  
 point in this respect is the assumed NS mass spectrum since the cooling evolution 53  
 is very sensitive to the star mass. However, as it is shown, new results largely agree 54  
 with previous ones and offer further support to the idea that the Gould Belt is the 55  
 nursery of the local NS population. 56

## 2. Model 57

The details of our model have been presented in Paper I; its main features are 58  
 summarized below. We assume that NSs are continuously born in the galactic disc 59  
 (up to a distance of  $\sim 3$  kpc from the Sun) and in the Gould Belt at a constant rate (see 60  
 Figure 1). The rates are different in the disc and in the Belt and have been estimated 61  
 from available SN progenitors counts (Tammann et al., 1994; Grenier, 2000). Both 62  
 the spatial and the cooling history of newborn NSs is then followed as they evolve 63



*Figure 1.* A sketch of the initial spatial distribution. It is a projection to the plane perpendicular to the galactic one. Stars are born in the Gould Belt, which is inclined to the galactic plane by 18 degrees, and in the galactic disc. Star producing regions are shown with thick lines.

64 in the galactic potential. Typically we calculate  $\sim 10^4$  evolutionary tracks and then  
 65 normalize our results to the actual number of NSs born in the considered volume  
 66 ( $\sim 1000$  NSs in a sphere of radius 3 kpc centered on the Sun) during a 4.25 Myrs  
 67 time interval. NSs cooling curves by Kaminker et al. (2002) have been used to derive  
 68 the NS temperature at each time step. The duration of the calculation is fixed by the  
 69 request that the surface temperature of the lightest (i.e. the hottest) NSs is higher than  
 70  $10^5$  K. Cooler NSs could not have been detected by ROSAT even if they are as close  
 71 as 10 pc. Since young cooling NSs are expected to emit most of their luminosity  
 72 at UV/soft X-ray energies ( $\sim 20$ – $200$  eV, corresponding to temperatures  $\sim 10^5$ – $10^6$   
 73 K), interstellar absorption must be accounted for. The PSPC count rate is finally  
 74 obtained from the unabsorbed flux, which corresponds to the given temperature,  
 75 radius and distance of the star, and from the value of the column density. This allows  
 76 us to construct the  $\log N - \log S$  curve for close-by, cooling NSs.

77 Results presented in Paper I rely on a particular choice for a set of free parameters  
 78 which enter the model. They reflect our incomplete knowledge of some properties  
 79 of the NS population, and are mainly related to: (i) the spatial distribution of NS  
 80 progenitors; (ii) the NS mass distribution; (iii) the NS kick velocity distribution;  
 81 (iv) the NS emitted spectrum. In Paper I we assumed that NSs are uniformly born  
 82 in the Belt, modeled as a thin disc 500 pc in radius, and that their initial velocity  
 83 distribution is represented by a single maxwellian with a mean velocity of 225 km  
 84  $s^{-1}$ . The mass spectrum was taken to be flat in the mass range  $1.1 M_{\odot} \leq M \leq$   
 85  $1.8 M_{\odot}$ . Cooling NSs were assumed to emit a pure blackbody spectrum without  
 86 allowance for possible deviations arising from reprocessing in an atmosphere and/or  
 87 by a reduce surface emissivity.

88 Here we address all these points in more detail. In particular we assess the  
 89 effects of relaxing the assumptions of Paper I on the computed  $\log N - \log S$  curve.  
 90 The main changes are described below. As it will be shown in the next section,  
 91 the original results presented in Paper I are not much influenced. With respect to  
 92 Paper I we introduce four modifications.

- 93 – We use a slightly different spatial distribution of NS progenitors taking the Gould  
 94 Belt radius to be 300 pc (instead of 500 pc, see Figure 1). The total birthrate in  
 95 the Belt was the same.
- 96 – Natal kicks were drawn from the complete velocity distribution of Arzoumanian  
 97 et al. (2002), described by two maxwellians with total average velocity  
 98  $\sim 540$  km  $s^{-1}$ , instead of a single maxwellian with average velocity  $\sim 225$  km  $s^{-1}$ .
- 99 – We account for the possible reduced emissivity of the star surface, as suggested  
 100 by the case of RX J1856.5-3754 (e.g. Drake et al., 2002). This has been mimicked  
 101 using a radiation radius  $R_{\text{rad}} \sim 0.32R$ , so that  $L = 4\pi R_{\text{rad}}^2 \sigma T_{\text{eff}}^4 \sim 0.1 L_{\text{BB},R}$ .<sup>1</sup>

<sup>1</sup>Such a strong reduction of emissivity of course influences cooling curves. However, we do not attempt to take into account self-consistently in this paper. A self-consistent calculations will be presented separately. Also possible influence of an atmosphere should be taken into account.

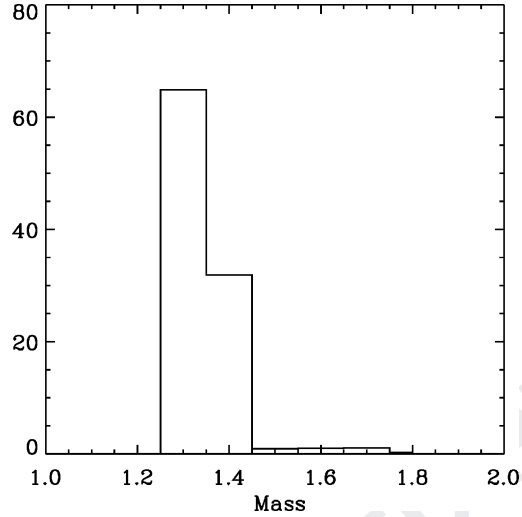


Figure 2. Mass distribution for young close-by NSs. Stars were distributed in eight bins from 1.1 to 1.8 solar masses. The vertical axis shows percentage in each bin.

- A more realistic mass spectrum of NSs, peaked around 1.3–1.4  $M_{\odot}$  (see 102  
Figure 2), has been derived and incorporated in the simulations instead of the 103  
flat one used before. 104

This last point requires some more comments. Performing a population synthesis 105  
of cooling NSs demands for the NS mass spectrum, since cooling curves depend 106  
on mass (e.g. Kaminker et al., 2002). As noted by Woosley et al. (2002), at present 107  
models do not allow a precise determination of the NS mass spectrum. However, 108  
given the dependence of the cooling curves on mass (see below), even a rough 109  
estimate is enough for the case at hand. In our calculations we use cooling curves 110  
for NS masses in the range  $1.1 M_{\odot} < M < 1.8 M_{\odot}$ ; masses are grouped in eight 111  
bins. Cooling models show that there is a critical value for the mass ( $\sim 1.35 M_{\odot}$  in 112  
the case of Kaminker et al. (2002), the exact value depends on model assumptions) 113  
across which the cooling history significantly changes. NSs with masses below the 114  
critical value have similar cooling histories and remain hot for a relatively long time 115  
( $T = 10^5$  K after 4.25 Myrs). Intermediate mass stars ( $\sim 1.4$ – $1.5 M_{\odot}$ ) cool down to 116  
 $10^5$  K in about the same time but have lower temperatures during the first million 117  
years of their evolution in comparison with less massive stars. NSs with masses 118  
 $M > 1.5 M_{\odot}$  experience much faster cooling and become completely invisible (at 119  
X-ray energies) in a few hundred thousand years or even earlier. 120

To our end, the most important point is estimating the number of NSs above 121  
and below the critical mass. In our discrete description this amounts to assess the 122  
number of stars in the first three mass bins relative to remaining five. In order to 123  
do so, we proceed as follows. At first we take massive stars closer than 500 pc 124  
(i.e. with known parallax  $> 0.002$ ) from the Hipparcos catalog ESA, 1997. Stars 125

126 from B2 to O8 are considered here to be NS progenitors. Spectral classes presented  
 127 in the catalog are then transformed into masses, although we are aware that this is a  
 128 very rough procedure. NS masses are finally obtained from the model described in  
 129 Timmes et al. (1996) and Woosley et al. (2002). To do it we use a fit of Figure 14 in  
 130 Woosley et al. (2002) and assumed that all stars less massive than  $\sim 11 M_{\odot}$  produce  
 131 NSs of the same mass, i.e.  $1.27 M_{\odot}$ . In all other cases the baryonic NS mass is  
 132 calculated from the mass of the progenitor according to

$$M_{\text{bar}} = \begin{cases} 0.067 M + 0.567 & 11 M_{\odot} < M < 15 M_{\odot} \\ \text{const} = 1.567 & 15 M_{\odot} \leq M \leq 20 M_{\odot} \\ 0.0867 M - 0.167 & M > 20 M_{\odot}. \end{cases} \quad (1)$$

133

134 The NS gravitational mass (which is used in our calculation) is calculated ac-  
 135 cording to  $M_{\text{bar}} - M_{\text{grav}} = 0.075 M_{\text{grav}}^2$  (Timmes et al., 1996), here and in the formula  
 136 below masses are in the solar units. All stars from the solar proximity contributed to  
 137 the final distribution with some coefficient, inversely proportional to their lifetime  
 138 ( $\log t = 9.9 - 3.8 \log M + \log^2 M$ ).

139 Within the 500 pc sphere the number of progenitors with  $M < 13.85 M_{\odot}$   
 140 (which give a  $1.35 M_{\odot}$  NS) is about twice higher than expected from the Salpeter  
 141 mass function. Such an enhancement is mostly connected with the Gould Belt. We  
 142 find that about 2/3 of NSs have mass  $< 1.35 M_{\odot}$  and that most NSs fall into the  
 143 1.3 and  $1.4 M_{\odot}$  bins (see Figure 2). The contribution of massive ( $> 1.5 M_{\odot}$ ) NSs is  
 144 negligible (about 3% by number). This is a special feature of the solar proximity.

145 This mass spectrum is in reasonable correspondence with mass determinations  
 146 in binary radio and X-ray pulsars. We note, that the peak at  $1.3 M_{\odot}$  is due to the  
 147 assumption (see Timmes et al., 1996) that all NSs below  $\sim 11 M_{\odot}$  produce NSs  
 148 of nearly the same mass,  $\sim 1.27 M_{\odot}$ . However, smearing the peak over the first  
 149 three mass bins would produce about the same results for  $\log N - \log S$  since the  
 150 cooling curves for these masses are very similar in the time interval of interest for  
 151 our calculations (see Kaminker et al., 2002). Although it represents just a rough  
 152 estimate, this spectrum is better, in our opinion, than the flat one we used before,  
 153 being closer to the one obtained from the mass determinations in binary systems  
 154 (mostly in binary radio pulsars), and in general it is in better correspondence with  
 155 expectations about young NSs in the solar vicinity.

156

### 3. Results

#### 157 3.1. LOG $N$ –LOG $S$ CURVES

158 The  $\log N - \log S$  distributions for young cooling INSs originated from the Gould  
 159 Belt and circumsolar parts of the galactic disc have been calculated for different  
 160 parameters characterizing the NSs mass, velocity, and spatial distributions. The

final variant, shown in Figure 5, includes all the four modifications discussed in the previous section. For comparison observational points representing the  $\log N - \log S$  distribution of isolated close-by NSs are also shown. These sources include three small groups of INSs: the seven ROSAT radioquiet INSs (“the Magnificent seven”), four young close-by radio pulsars with detected thermal radiation (“the three musketeers” and PSR B1929 + 10) and Geminga together with the Geminga-like object 3EG J1835 + 5918 (see Paper I for details). The “Magnificent seven” is characterized by absence of detected radio emission and long spin periods (about 5–10 s). Geminga and 3EG J1835 + 5918 are known  $\gamma$ -ray sources.

Symbols for the observed data are in correspondence with the type of the faintest object which contributes at a given flux (filled: ROSAT radioquiet INSs, empty: other sources). That is if the faintest object at a given flux is one of the magnificent seven then we plot a filled diamond, if not – an empty circle. Error bars represent poissonian errors. The two limits on the number of INSs in the ROSAT data obtained by Rutledge et al. (2003) (labeled BSC) and Schwope et al. (1999) (labeled RBS) are also shown.

Let us discuss first the relative effect of each of the four modifications mentioned above. Clearly, a reduced surface emissivity and higher average kick velocity act in decreasing the number of observable sources at a given flux, while a higher fraction of low mass (i.e. hotter) NSs and a more compact initial distribution tend to increase it.

In Figure 3 we compare our previous result for the disc alone and disc + Belt with the new ones for disc + Belt. To obtain these new curves we considered either a smaller radiation radius and flat mass spectrum (dot-dot-dashed curve) or the new mass spectrum together with standard emissivity (dot-dashed curve). One can see that the reduced emissivity and the new peaked mass spectrum move the  $\log N - \log S$  curve (down and up, respectively) by nearly half order of magnitude each, with a net combined effect of slightly decreasing the number of observable sources. Here the spatial distribution of NS progenitors and kick velocities are the same as in the original calculations. From the next figure (Figure 4) it is apparent that the same effect is produced when a more compact initial distribution and higher kick velocities are introduced, keeping the original assumptions about the star emissivity and mass spectrum. Together these two modifications tend to slightly increase the number of observable sources.

Our final results for  $\log N - \log S$  are presented in Figure 5. Here all four effects are taken into account. The general conclusion is that modifications do not change our results significantly. Our estimate is well below the BSC limit by Rutledge et al. (2003) and in correspondence with the RBS limit (Schwope et al., 1999).

### 3.2. SPATIAL DISTRIBUTION

In this section we briefly discuss the spatial distribution of observable INSs. Despite the main focus of our work has been of the  $\log N - \log S$  curve, the present distribu-

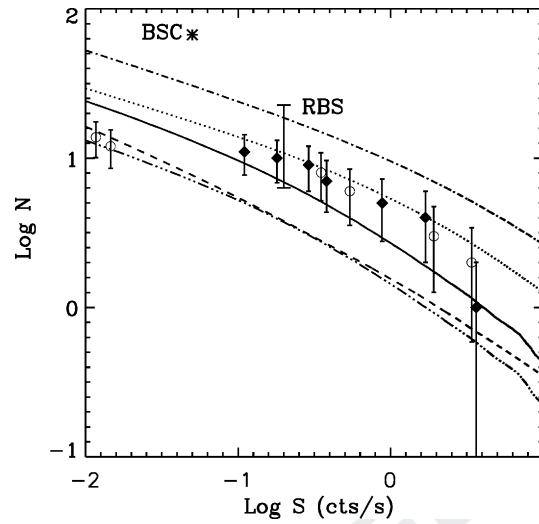


Figure 3. Relative contribution of the “atmospheric effect” and mass spectrum are shown. Dotted line – the “old” model for disc + Belt contribution. Dashed line – the “old” model without a Gould Belt. Dot-dot-dashed (the lowest line) – “atmospheric effect.” Dot-dashed – effect of the new mass spectrum. Solid line – both effects together. Symbols which show observational points (filled diamonds or open circles) are in correspondence with the type of the faintest object (a ROSAT INS or not) which contributes to the total number at the specified count rate (see text for more details). RBS and BSC are limits on the number of bright INNs in ROSAT data obtained by Schwöpe et al. (1999) and Rutledge et al. (2003), respectively.

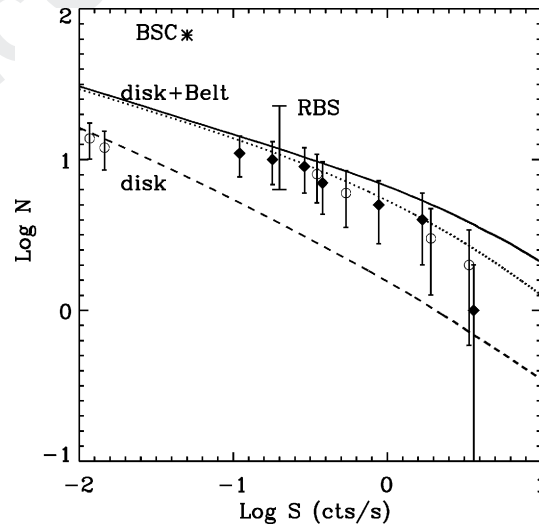


Figure 4. Contributions of variations of initial spatial distribution and kick velocity distribution are shown. Solid line:  $R_{\text{Belt}} = 300$  pc and kick velocity distribution as in Arzoumanian et al. (2002).

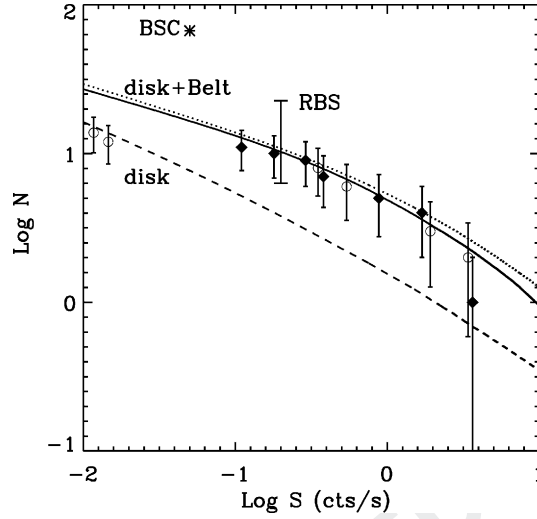


Figure 5. A “new” model (solid line:  $R_{\text{Belt}} = 300$  pc, new mass spectrum, an “atmospheric effect”, new kick velocity distribution) in comparison with older results (dotted line – disc + Belt, dashed line – only disc).

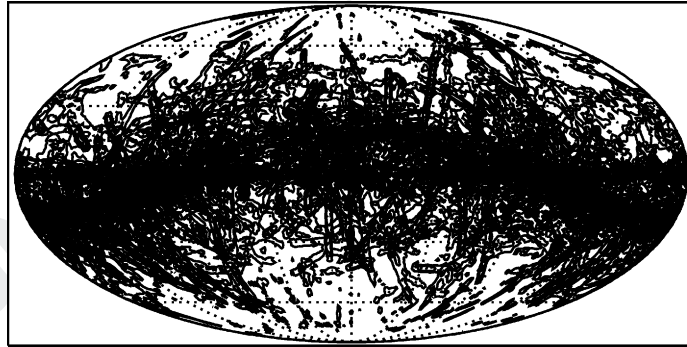


Figure 6. Projected distribution of cooling INSs in the sky in galactic coordinates. Only sources with count rate  $>0.05$   $\text{cts s}^{-1}$  are accounted for. The plot shows contours of constant INS number density per square degree. Darker areas close to the Belt or/and to the galactic plane correspond to  $\sim 0.001$  sources/square degree. The presence of the Belt produces a tilt in the higher projected density region which is visible in the figure.

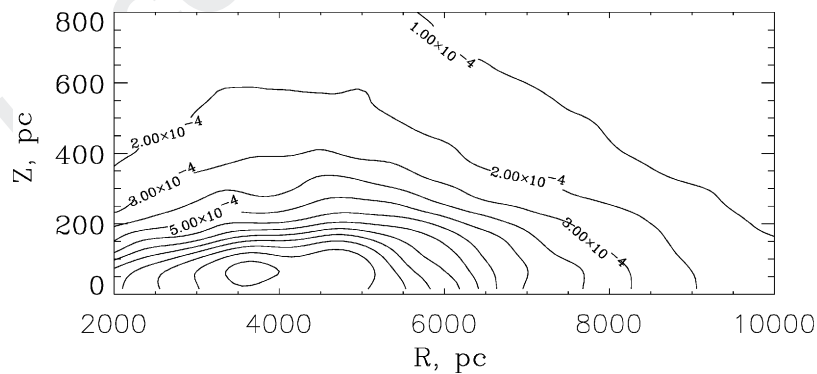
tion of cooling INSs on the sky is obtained as a by-product of our evolutionary code. 202  
 For illustrative purposes we report in Figure 6 the projected spatial distribution of 203  
 relatively bright coolers (count rate  $> 0.05$  PSPC  $\text{cts s}^{-1}$ ) in galactic coordinates. 204  
 In the figure we present contours of constant number of INS. The galactic plane and 205  
 the Gould Belt are clearly visible. High-latitude regions are nearly free of young 206  
 NSs. 207



208 This result should be taken with care as far as we used a simplified initial  
 209 spatial distribution for the progenitors and did not take into account any detailed  
 210 ISM structure around the Sun. The figure just illustrates some general features  
 211 of the distribution. The two main features of our model are apparent from the  
 212 plot of Figure 6 which shows that the highest projected density of sources is close  
 213 to the galactic plane. The presence of the Belt produces the tilt of the high den-  
 214 sity region towards low galactic latitudes. All small scale details are due to in-  
 215 dividual tracks of calculated INSs and should not be treated as predictive of any  
 216 “fine structure.”

217 As expected, sources are strongly concentrated towards the galactic plane and  
 218 the Gould Belt. Only about 12% of sources with ROSAT count rate  $>0.01$  cts  
 219  $s^{-1}$  are found at  $|b| > 40^\circ$ . About 20% of sources lie outside the belt  $\pm 30^\circ$  from  
 220 the galactic plane, while  $\sim 50\%$  are expected to be within  $\pm 12^\circ$  from the plane of  
 221 the Galaxy (brighter sources are more strongly concentrated towards the galactic  
 222 plane and the Gould Belt since they correspond to younger INSs). Although the very  
 223 strong concentration towards the galactic plane may reflect the assumption that NSs  
 224 are born exactly in the (infinitesimally thin) galactic disc, the source distribution at  
 225 higher latitudes should be real.

226 Finally we would like to stress that the distribution of young NSs around the  
 227 Sun is definitely different from the full NS spatial distribution, which is dominated  
 228 by old stars (age  $> 10^7$  years). For comparison the latter is shown in Figures 7 and 8  
 229 for two different assumptions about NS formation rate distribution.



*Figure 7.* Distribution of all isolated NSs in the Galaxy in the  $R$ - $z$  plane. The data is calculated by a Monte Carlo of  $>10,000$  individual tracks on a fine grid (10 pc in  $z$  direction and 100 pc in  $R$  direction). Curves were smoothed, all irregularities are of statistical nature. Kick velocity is assumed following Arzoumanian et al. (2002). NSs are born in the thin disc with semithickness 75 pc. No NS born inside  $R = 2$  kpc and outside  $R = 16$  kpc are taken into account. NS formation rate is assumed to be proportional to the square of the ISM density at the birthplace. Results are normalized to have in total  $5 \times 10^8$  NSs born in the described region. Density contours are shown with a step  $0.0001 \text{ pc}^{-3}$ .

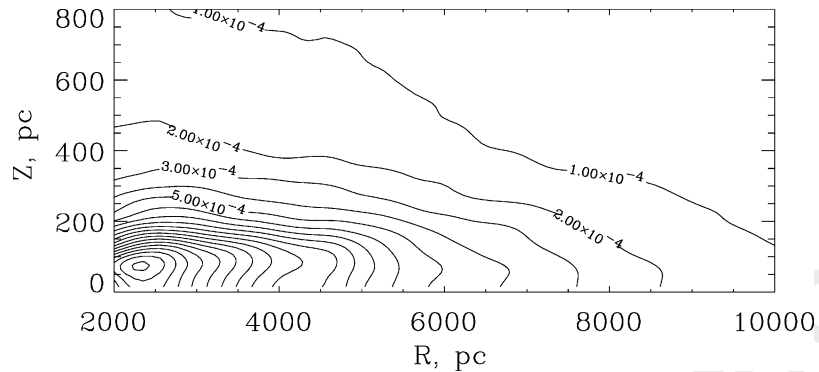


Figure 8. Distribution of all isolated NSs in the Galaxy in the  $R$ - $z$  plane. All parameters are as in the previous figure except the distribution of NS formation rate, it is assumed to be proportional to  $[\exp(-z/75 \text{ pc}) \exp(-R/4 \text{ kpc})]$ . Curves were smoothed as in the previous picture. It is clearly seen that in that case NSs are stronger concentrated towards the galactic center, then in the case of NS formation rate proportional to the square of the ISM density.

We do not expect new identifications of bright ( $>0.1 \text{ cts s}^{-1}$ ) cooling INSs at large galactic latitudes. Most of the unidentified ROSAT objects (still there should be tens of them for count rates  $>0.01 \text{ cts s}^{-1}$ ) should be in crowded fields at  $\pm 30^\circ$  from the plane of the Milky Way. Some can be identified as EGRET sources among which young INSs from the Gould Belt should be numerous (Grenier, 2000).

#### 4. Conclusions

In this paper we presented results of more advanced population synthesis model (in comparison with Paper I) of young close-by INSs. Account of new effects does not change our previous conclusion that observed cooling INSs in the solar vicinity originated in the Gould Belt.

We also modeled distribution of bright cooling close INSs on the sky. Most of them are situated close to the galactic plane and to the plane of the Gould Belt.

More detailed calculation is necessary. It is especially important to take into account effects of the most outer layers of NSs on the spectrum and detailed distribution of NS progenitors. After parameters of a population synthesis model are fixed it is possible to use  $\log N - \log S$  distribution as an independent test of models of thermal evolution of NSs.

#### Acknowledgements

We want to thank Dmitry Yakovlev and his colleagues for putting their cooling model to our disposal, and Hovik Grigorian for very useful comments. SP also

250 thanks Isabelle Grenier, David Blaschke and Dima Voskresensky for discussions.  
 251 The work of MP and SP was partly supported by the RFBR grant 03-02-16068 and  
 252 by the program “Universities of Russia” grant 02.03.013/2. The work of MC, SP,  
 253 AT and RT was partially supported by the Italian Ministry for Education, University  
 254 and Research (MIUR) under grant COFIN-2000-MM02C71842.

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**Queries**

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